VECTOR POTENTIAL OF THE BASIS IN THE SPHERICAL

COORDINATES

For Electromagnetism A, Univ. Tokyo (2016)

Descriptions: Correction to the last question in QI, in the 2016 final exam, which is about the vector potential for the basis in the spherical coordinates

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1 Divergence of the spherical-coordinate basis

In the three-dimensional spherical coordinates, the basis are given

$$\boldsymbol{e}_{\mathrm{r}} = \begin{pmatrix} \cos\phi\sin\theta\\ \sin\phi\sin\theta\\ \cos\theta \end{pmatrix}, \quad \boldsymbol{e}_{\theta} = \begin{pmatrix} \cos\phi\cos\theta\\ \sin\phi\cos\theta\\ -\sin\theta \end{pmatrix}, \quad \boldsymbol{e}_{\phi} = \begin{pmatrix} -\sin\phi\\ \cos\phi\\ 0 \end{pmatrix}. \tag{1}$$

Only the last basis vector, that is, e_{ϕ} , gives the vanishing divergence,

$$\boldsymbol{\nabla} \cdot \boldsymbol{e}_{\phi} = \boldsymbol{0}. \tag{2}$$

According to a vector identity, div curl = 0, as seen in the Gauss law for the magnetic field, the basis vector e_{ϕ} can be expressed as $e_{\phi} = \nabla \times A$. Let us find this vector-potential field A in below.

2 Expression of the derivative operator ∇ in the spherical coordinates

Let us present the expression first.

$$\boldsymbol{\nabla} = \boldsymbol{e}_{\mathrm{r}} \frac{\partial}{\partial \mathrm{r}} + \boldsymbol{e}_{\theta} \frac{1}{\mathrm{r}} \frac{\partial}{\partial \theta} + \boldsymbol{e}_{\phi} \frac{1}{\mathrm{r} \sin \theta} \frac{\partial}{\partial \phi}.$$
(3)

We then express the vector potential as

$$\mathbf{A}(\mathbf{r}) = \mathbf{A}_{\mathbf{r}} \mathbf{e}_{\mathbf{r}} + \mathbf{A}_{\theta} \mathbf{e}_{\theta} + \mathbf{A}_{\phi} \mathbf{e}_{\phi}.$$
 (4)

Finally, we need to perform the following

$$\boldsymbol{\nabla} \times \mathbf{A} = \left(\boldsymbol{e}_{\mathrm{r}} \frac{\partial}{\partial \mathrm{r}} + \boldsymbol{e}_{\theta} \frac{1}{\mathrm{r}} \frac{\partial}{\partial \theta} + \boldsymbol{e}_{\phi} \frac{1}{\mathrm{r} \sin \theta} \frac{\partial}{\partial \phi} \right) \times \left(A_{\mathrm{r}} \boldsymbol{e}_{\mathrm{r}} + A_{\theta} \boldsymbol{e}_{\theta} + A_{\phi} \boldsymbol{e}_{\phi} \right)$$
(5)

To carry out this calculation, we may make use of the following fomrula,

$$(\mathbf{A}\partial) \times (\mathbf{f}\mathbf{B}) = (\partial \mathbf{f})(\mathbf{A} \times \mathbf{B}) + \mathbf{f}(\mathbf{A} \times \partial \mathbf{B}), \tag{6}$$

which has been derived when we considered the expression of curls in the cylindrical coordinates.

3 Curl in the spherical coordinates

The obtained result reads

$$\nabla \times \mathbf{A} = \mathbf{e}_{\mathrm{r}} \left(\frac{1}{\mathrm{r}} \frac{\partial A_{\phi}}{\partial \theta} - \frac{1}{\mathrm{r} \sin \theta} \frac{\partial A_{\theta}}{\partial \phi} + A_{\phi} \frac{\cos \theta}{\mathrm{r} \sin \theta} \right) + \mathbf{e}_{\theta} \left(-\frac{\partial A_{\phi}}{\partial \mathrm{r}} + \frac{1}{\mathrm{r} \sin \theta} \frac{\partial A_{\mathrm{r}}}{\partial \phi} - \frac{A_{\phi}}{\mathrm{r}} \right) + \mathbf{e}_{\phi} \left(\frac{\partial A_{\theta}}{\partial \mathrm{r}} - \frac{1}{\mathrm{r}} \frac{\partial A_{\mathrm{r}}}{\partial \theta} + \frac{A_{\theta}}{\mathrm{r}} \right).$$
(7)

The condition we reqire is

$$\boldsymbol{e}_{\boldsymbol{\varphi}} = \boldsymbol{\nabla} \times \boldsymbol{A}, \tag{8}$$

so that we need to demand

$$\frac{1}{r}\frac{\partial A_{\phi}}{\partial \theta} - \frac{1}{r\sin\theta}\frac{\partial A_{\theta}}{\partial \phi} + A_{\phi}\frac{\cos\theta}{r\sin\theta} = 0, \qquad (9)$$

$$-\frac{\partial A_{\phi}}{\partial r} + \frac{1}{r \sin \theta} \frac{\partial A_{r}}{\partial \phi} - \frac{A_{\phi}}{r} = 0, \qquad (10)$$

$$\frac{\partial A_{\theta}}{\partial r} - \frac{1}{r} \frac{\partial A_{r}}{\partial \theta} + \frac{A_{\theta}}{r} = 1.$$
(11)

to be satisfied.

4 A special case: $A(r) = A_r(r)e_r$

In the exam, I asked to find a special solution for the equations above, which is the case when the vector potential is parallel to the radial basis e_r . That is,

$$\mathbf{A}(\mathbf{r}) = A_{\mathrm{r}}(\mathbf{r}, \theta, \phi) \boldsymbol{e}_{\mathrm{r}},\tag{12}$$

which also means $A_{\theta} = A_{\phi} = 0$. The first equation of the three differential equations above is satisfied in a trival way. The second and third equations turn to

$$\frac{\partial A_{\rm r}}{\partial \phi} = 0, \tag{13}$$

$$\frac{\partial A_{\rm r}}{\partial \theta} = -r. \tag{14}$$

The first equation means that A_r does not depend on variable ϕ , so that $A_r = A_r(r, \theta)$. Inserting this result to the second equation gives rise to

$$A_{r}(r,\theta) = -r\theta + X(r), \qquad (15)$$

where $X(\mathbf{r})$ is an arbitrary function of \mathbf{r} . Hence,

$$\boldsymbol{e}_{\phi} = \boldsymbol{\nabla} \times \{(-\mathbf{r}\boldsymbol{\theta} + \mathbf{X}(\mathbf{r})) \, \boldsymbol{e}_{\mathbf{r}}\}. \tag{16}$$